

# 2D Amorphous/Crystalline $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$ Nanosheet Heterostructures with Improved Capability for $\text{H}_2$ and $\text{NO}_2$ Sensing

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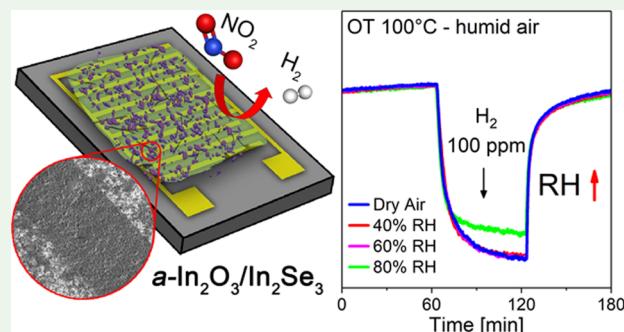
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**ABSTRACT:** Spontaneous degradation of 2D transition-metal dichalcogenides/chalcogenides (TMDs/MCs) gas sensors in dry/wet air represents one of the most significant drawback of these interfaces, hampering the reproducibility of the baseline resistance and sensor's signal stability (i.e., sensor's creep). Herein, we report a simple protection strategy stimulating the formation of a self-assembled oxide ( $a\text{-MO}_x$ ) over TMDs/MCs, which promotes effective passivation of the underlying surface and excellent gas sensing response. Liquid-phase-exfoliated few-layers 2D- $\text{In}_2\text{Se}_3$  have been annealed in air at 180 °C for 24 h to yield an  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure comprising a self-assembled  $a\text{-In}_2\text{O}_3$  amorphous skin (5–10 nm) over 2D-crystalline  $\text{In}_2\text{Se}_3$  (5–30 nm). The isomorphic conversion of  $\text{In}_2\text{Se}_3$  into  $a\text{-In}_2\text{O}_3$  specifically enables the layered shape of the precursor 2D- $\text{In}_2\text{Se}_3$  to be preserved after annealing, therefore providing all the surface-to-volume advantages of 2D interfaces. The excellent baseline and sensor's signal reproducibility to  $\text{H}_2$  (5–100 ppm) and  $\text{NO}_2$  (400 ppb–1 ppm) after 1 year of delivery at 100 °C operating temperature demonstrated that the oxide skin effectively passivates the underlying 2D- $\text{In}_2\text{Se}_3$  from further oxidation. Significantly, the  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure shows better  $\text{H}_2$  sensing response with respect to 2D TMDs/MCs sensors, with experimental detection limits as low as 5 ppm  $\text{H}_2$  and 400 ppb  $\text{NO}_2$ , with associated RR ( $R_a/R_g$ ) = 2.1 to 100 ppm  $\text{H}_2$  and RR ( $R_g/R_a$ ) = 2.3 to 1 ppm  $\text{NO}_2$  in dry air. A charge carrier mechanism between the  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure and  $\text{H}_2$ ,  $\text{NO}_2$ , and  $\text{H}_2\text{O}$  molecules is presented to discuss the humidity cross response to  $\text{H}_2$  and  $\text{NO}_2$ . The passivation strategy here proposed can be extended to a large variety of TMDs/MCs, opening new perspectives for the effective exploitation of layered amorphous gas-sensing interfaces.

**KEYWORDS:**  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure, amorphous  $\text{In}_2\text{O}_3$ , ( $a\text{-In}_2\text{O}_3$ ), chemoresistive,  $\text{H}_2$  sensor,  $\text{NO}_2$  sensor



## 1. INTRODUCTION

Hydrogen as a clean and renewable fuel is one of the most promising energy carriers for the reduction of the environmental impact of the automotive and transport industry. Despite that, hydrogen-related hazards in production, storage, and transport (explosion range 4–75% vol  $\text{H}_2$  in air) represent crucial issues for its safe exploitation, challenging new sensor platforms suitable to detect colorless and odorless hydrogen leaks, down to a few parts per million (ppm). Nitrogen dioxide, on the other hand, is a ubiquitous gas deriving from emissions from vehicles, power plants, and factories, mostly as a byproduct of the combustion process. EPA's standards for  $\text{NO}_2$  limit people's short-term exposures to 100 ppb, making its actual measure challenging in complex urban and indoor gaseous atmospheres.<sup>1</sup>

Chemoresistive  $\text{H}_2$  and  $\text{NO}_2$  sensors based on different metal oxides (MO) have been recently reviewed,<sup>2,3</sup> highlighting the excellent  $\text{H}_2/\text{NO}_2$  sensing properties of different sensor architectures at operating temperatures (OTs) in the range of 200–260 °C. More recently, new  $\text{H}_2/\text{NO}_2$  gas platforms, based on few layers of exfoliated van der Waals

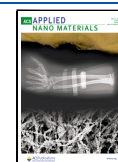
(vdW) transition-metal dichalcogenides (TMDs) and metal chalcogenides (MCs), operating at lower temperatures (i.e., 25–100 °C), have been investigated as possible alternatives to traditional  $\text{H}_2/\text{NO}_2$  MO sensors.<sup>4</sup> These TMDs/MCs nanosheet interfaces, which can be easily engineered into complex 0D–1D–3D architectures<sup>5</sup> and 2D/2D heterostructures,<sup>6–8</sup> provide all the surface-to-volume advantages of layered interfaces, enabling the transduction of sub-ppm concentrations of  $\text{NO}_2$  and ppm level of less electronegative gases like  $\text{H}_2$  and CO.

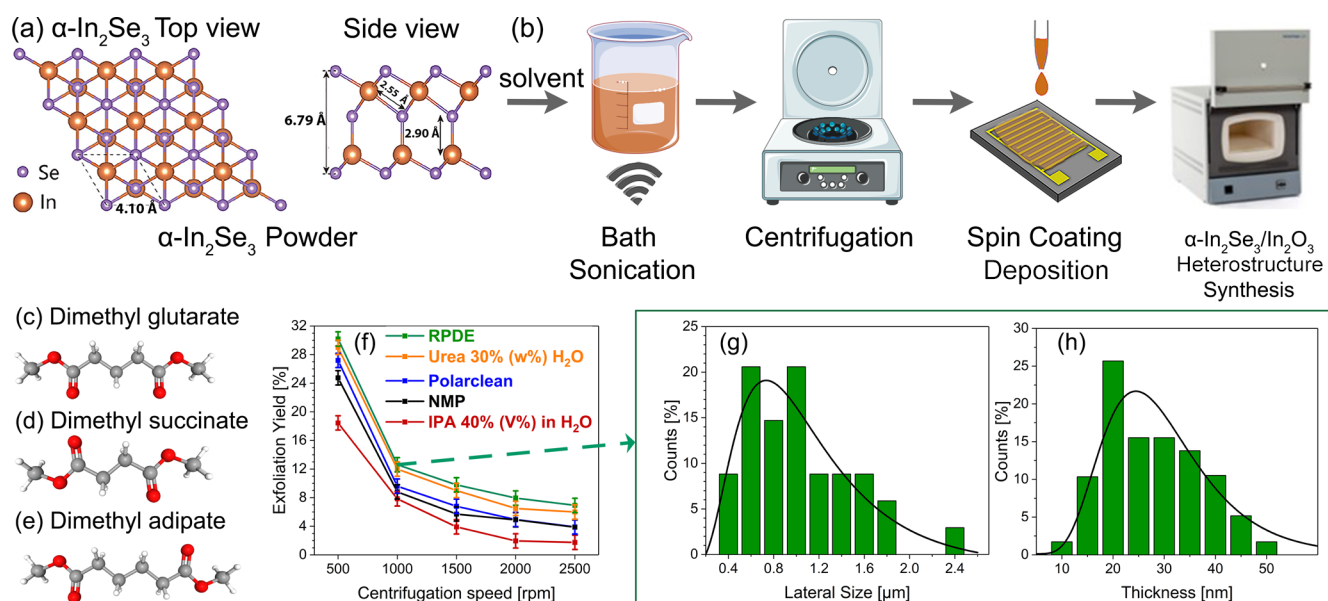
Remarkably, few layers of exfoliated TMDs and MCs chemoresistive sensors are subjected to fast oxidation/degradation in dry/humid laboratory conditions.<sup>9</sup> Their thermodynamic instability, demonstrated by the spontaneous

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**Figure 1.** (a) Planar and side views of the Se–In–Se–In–Se quintuple layer blocks of the crystal structure of the  $\alpha$ - $\text{In}_2\text{Se}_3$  phase. Reproduced from ref 21: Nian, T.; Wang, Z.; Dong, B. Thermoelectric Properties of  $\alpha$ - $\text{In}_2\text{Se}_3$  monolayer. *Appl. Phys. Lett.* **2021**, *118*(3). Copyright 2021 AIP Publishing; (b) schematization of the sonication-assisted LPE process comprising spin-coating deposition of the flakes over interdigitated patterns and thermal treatment in air at 180 °C for 24 h to synthesize the  $\alpha$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure; (c–e) ball-and-stick representation of the Rhodiasolv RPDE solvent's main constituents (red: oxygen, gray: carbon, and white: hydrogen); (f) exfoliation yield (%) as a function of the centrifugation speed after 15 min for different solvents (RPDE, PolarClean, NMP, urea 30% (w%) in water, and 40% IPA in water). (g,h) SEM and AFM statistical analysis of the lateral size and thickness distribution of RPDE exfoliated  $\text{In}_2\text{Se}_3$  flakes corresponding to 1000 rpm and 15 min centrifugation.

substitution of the chalcogen atoms (S, Se, and Te) operated by ambient  $\text{O}_2$  ( $\Delta G < 0$ ), is an inexplicably overlooked issue when exploiting new TMDs and MCs platform interfaces for electronic applications. Specifically, the surface degradation process, eventually accelerated at high OTs, is responsible, in the long run, for irreversible drift phenomena of the electrical response (i.e., sensor signal creep), limiting the exploitation of these interfaces.

In the case of binary III–IV group 2D-layered semiconductor selenides, the spontaneous oxidation process initiates with the formation of an amorphous skin of  $\alpha$ - $\text{MO}_x$  ( $M = \text{In, Ga, Sn, Ge}$ ) advancing from the surface toward the inner layers as reported for  $\text{InSe}$ ,<sup>10</sup>  $\text{GaSe}$ ,<sup>11</sup>  $\text{GeSe}$ ,<sup>12</sup> and  $\text{SnSe}$ ,<sup>13</sup> addressing new strategies to passivate the sensor surface from fast degradation.

We have recently reported<sup>14,15</sup> a successful synthesis, which passivates the nanoplate's surface, yielding the formation of either a protective amorphous metal-oxide skin layer of  $\alpha$ - $\text{SnO}_2$  over 2D- $\text{SnSe}_2$  to produce  $\alpha$ - $\text{SnO}_2/\text{SnSe}_2$  heterojunctions or fully  $\alpha$ - $\text{SnO}_2$  2D layered amorphous interfaces, with excellent long-term stability of the baseline resistance and excellent sensor' signal (i.e.,  $R_a/R_g$ ) reproducibility.

In this work, we demonstrate that the passivation/amorphization thermal oxidation process can be successfully extended to liquid-phase exfoliated commercial powders of binary III–IV group vdW 2D- $\text{In}_2\text{Se}_3$  semiconductor chalcogenides yielding  $\alpha$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructures, with unveiled  $\text{H}_2$  and  $\text{NO}_2$  responses in air/humid atmospheres.

We specifically report a liquid-phase exfoliation (LPE) procedure utilizing for the first time the Rhodiasolv RPDE eco-friendly commercial solvent to synthesize few layers of 2D- $\text{In}_2\text{Se}_3$  (20–30 nm thick) subsequently annealed in air at 180 °C for 24 h to promote the formation of the  $\alpha$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure. This paper, besides the effectiveness of the

passivation synthesis here validated, highlights that self-assembled 2D-amorphous/2D-crystalline heterostructures can detect oxidizing ( $\text{NO}_2$ ) and reducing ( $\text{H}_2$ ) gases in the same way as chemoresistive MO sensors, opening new perspectives to exploit the gas sensing properties of amorphous MO interfaces.

## 2. EXPERIMENTAL SECTION

**2.1. 2D  $\text{In}_2\text{Se}_3$  Exfoliation.**  $\text{In}_2\text{Se}_3$  crystals (Ossila—UK) were exfoliated in *N*-methyl-2-pyrrolidone (NMP, Sigma-Aldrich), Rhodiasolv RPDE (Solvay Novacare), Rhodiasolv PolarClean (Solvay Novacare), IPA (Sigma-Aldrich) 40 vol % in  $\text{H}_2\text{O}$  and Urea (Sigma-Aldrich) 30 wt % in  $\text{H}_2\text{O}$ , following a sonication-assisted LPE method comprising three steps:

- (i)  *$\text{In}_2\text{Se}_3$  exfoliation.* After 5 min of mortar grinding, 10 mg of  $\text{In}_2\text{Se}_3$  were dispersed in 20 mL of the chosen solvent inside a 50 mL Falcon tube and isothermally bath sonicated at  $24 \pm 2$  °C for 3 h (Elmasonic P60H at a 37 kHz frequency).
- (ii) *Solvent removal and purification.* After the first centrifugation (Neya 8 Bench Top centrifuge) at 5500 rpm for 15 min, two-thirds of the initial solvent were removed and replaced with ethanol (Sigma-Aldrich). After redispersion by sonication, three more rounds were carried out at 5000 rpm with subsequent replacement of the supernatant with fresh ethanol to ensure complete removal of the exfoliating medium.
- (iii) *Exfoliation yield.* To separate the exfoliated material, 15 min centrifugations were performed at different speeds (500, 1000, 1500, 2000, and 2500 rpm). Then, the 9/10 of the supernatant of each dispersion was taken by pipetting and dried overnight on a laboratory watch glass at 90 °C. Exfoliation yield [%] was calculated as the ratio between the initial and the final mass of the exfoliated dried material.

**2.2. Synthesis of the 2D  $\alpha$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  Heterostructure.** The 2D  $\alpha$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterostructure was synthesized by annealing the as-exfoliated nanosheets in static air at 180 °C for 24 h. After thermal

**Table 1. Surface Tensions ( $m[N/m]$ ) and Exfoliation Yields (%), Calculated as the Percentage Ratio of the Weight of the Flakes in the Final Dispersion with Respect to the Initial One, upon 1000 rpm and 15 min Centrifugation with Different Solvents, Utilizing  $WS_2$  and  $In_2Se_3$  Commercial Powders**

Solvent	Rhodiasolv RPDE	Rhodiasolv PolarClean	N-methyl-2-pyrrolidone (NMP)	IPA 40% (V %) $H_2O$	Urea 30% w $H_2O$
surface tension $\sigma_s$ $m[N/m]$	67 <sup>17</sup>	38 <sup>22</sup>	40 <sup>18</sup>	36 <sup>20</sup>	74 <sup>23</sup>
$In_2Se_3$ exfoliation yield [%]	13.0 $\pm$ 0.6	9.9 $\pm$ 0.7	8.6 $\pm$ 0.6	8.0 $\pm$ 0.7	12.8 $\pm$ 0.8
$WS_2$ exfoliation yield [%]	5.0 $\pm$ 0.6	16.2 $\pm$ 0.7	14.7 $\pm$ 0.6	13.4 $\pm$ 0.7	4.7 $\pm$ 0.8

treatment, flakes were redispersed in ethanol and deposited by spin coating over dedicated substrates for further characterization.

**2.3. Microstructural Characterization.** Exfoliated  $In_2Se_3$  microstructural features were investigated using SEM (Gemini FE-SEM 500), high-resolution transmission electron microscopy (HRTEM), and SAED at 200 kV by JEOL JEM 2010F (JEOL Ltd., Tokyo, Japan) and tapping-mode atomic force microscopy (AFM) (NT-MDT Ntegra, Spectrum). Thermogravimetric (TG–DTA) analyses were performed in air and  $N_2$  atmospheres using LINSEIS L81-I (Linseis GmbH, Germany). XPS spectra were collected using a PHI 1257 spectrometer with a monochromatic Al  $K\alpha$  source ( $h\nu = 1486.6$  eV) at a 250 meV spectral resolution. Grazing-incidence X-ray diffraction was performed using XRD-PAN Analytical X'PERT Pro with Cu  $K\alpha_1$  radiation ( $\lambda = 1.5406$  Å) with an incident angle of  $0.8^\circ$ . A PerkinElmer Lambda 1050 + UV/vis/NIR spectrophotometer (PerkinElmer Inc., Waltham, USA) equipped with a 100 mm integrating sphere was utilized for optical characterization. Reflectance spectra were collected in diffuse reflectance spectroscopy (DRS) configuration, and the Kubelka–Munk function was utilized to elaborate the signal.

**2.4. Electrical Characterization.** 30  $\mu$ L of the 2D  $a-In_2O_3/In_2Se_3$  nanosheets dispersion were spin-coated over  $Si_3N_4$  substrates provided with Pt finger-type electrodes (30  $\mu$ m apart) and backside heater placed inside a Teflon test cell (500  $cm^3$ ). Electrical resistance was measured (Agilent 34970A) at different OTs (25–150  $^\circ$ C range), as shown in Figure S1. Distinct gas concentrations were obtained by diluting certified mixtures (Nippon gases—IT) of  $NO_2$  (5 ppm in pure air) and  $H_2$  (250 ppm in pure air) with synthetic dry air utilizing mass flow controllers (MKS147), setting the total flow rate at 500 sccm/min. Different relative humidities (RH %) were obtained by mixing dry air with water-saturated air at 25  $^\circ$ C (i.e., RH % @ 25  $^\circ$ C). The RH % (RH % @ 25  $^\circ$ C) was measured before injecting the humid air into the test cell by using a thermo-hygrometer (HI9564, Hannah Instruments, Padua—IT). Exfoliated  $In_2Se_3$  gas sensing properties were assessed by using the following figures of merit: (i) baseline resistance (BLR): the sensor's resistance in air at equilibrium; (ii) sensor's signal: the relative resistances changes ( $R_{Gas}/R_{Air}$ ) and ( $R_{Air}/R_{Gas}$ ) to oxidizing/reducing gases; (iii) sensor sensitivity ( $S$ ): the slope of the calibration curve in the log–log plot, (iv) adsorption and desorption times ( $\tau_{ADS}$ ,  $\tau_{DES}$ ): the time needed to reach 90% of the sensors' signal steady-state value after gas adsorption or desorption.

### 3. RESULTS AND DISCUSSION

**3.1.  $In_2Se_3$  Material Exfoliation.** We exfoliated (Figure 1a) commercial crystals of  $\alpha-In_2Se_3$  by sonication-assisted LPE, collecting flakes with different lateral sizes and thicknesses depending on the centrifugation speed and time. The centrifuged suspension, after purification with ethanol, was spin-coated over interdigital electrodes and annealed in air at 180  $^\circ$ C for 24 h to promote the formation of the  $a-In_2O_3/In_2Se_3$  heterostructures. LPE exfoliation was carried out for the first time, utilizing Rhodiasolv RPDE, an eco-friendly and nontoxic commercial dibasic ester solvent,<sup>16</sup> with a chemical composition comprising three different dimethyl ester molecules (Figure 1c–e), as confirmed by Fourier transform infrared (FTIR) spectroscopy analysis shown in Figure S2 and product's data sheet specifications.<sup>17</sup> The RPDE solvent's

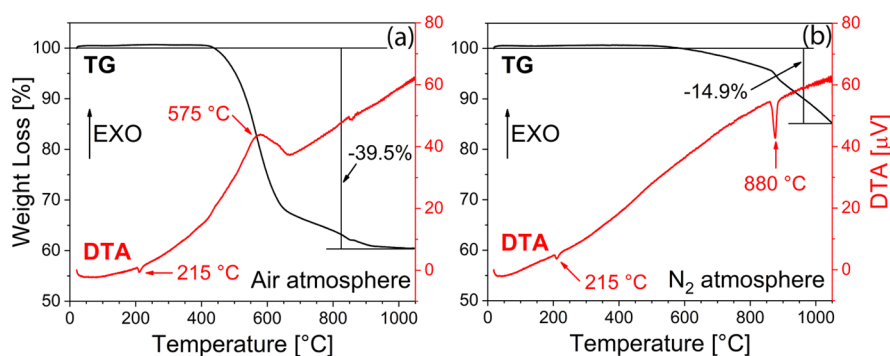
exfoliation yield (calculated for each centrifugation speed/time, as the percentage ratio of the weight of the flakes in the final dispersion with respect to the initial one), is shown in Figure 1f, with those obtained utilizing toxic NMP,<sup>18</sup> sustainable water soluble Rhodiasolv PolarClean,<sup>19</sup> 40 vol % IPA in water,<sup>20</sup> and Urea 30 wt % in water. The RPDE solvent exhibits better exfoliation yields (green line in Figure 1f) with respect to the others, demonstrating to be an excellent substitute for NMP and a suitable alternative to PolarClean. Finally, the centrifugation speed and time tune the morphology of the exfoliated  $In_2Se_3$ . Flakes' lateral size and thickness distributions, corresponding to a 1000 rpm centrifuge speed, are shown in Figures 1g,h and S3, with a calculated “aspect ratio” (AR) of  $\sim 40$  (AR = flakes lateral dimension/flakes thickness).

The exfoliation yields of different representative solvents, using  $In_2Se_3$  and  $WS_2$  as reference materials, upon 1000 rpm and 15 min centrifugation are compared in Table 1. RPDE and 30% (w %) urea/water are more effective in exfoliating  $In_2Se_3$  ( $\sim 13\%$ ) with respect to  $WS_2$  ( $\sim 5\%$ ). Conversely, PolarClean, NMP, and 40% (V %) IPA/water are more suitable to exfoliate  $WS_2$  ( $\sim 14\%$  average) than  $In_2Se_3$  ( $\sim 9\%$  average).

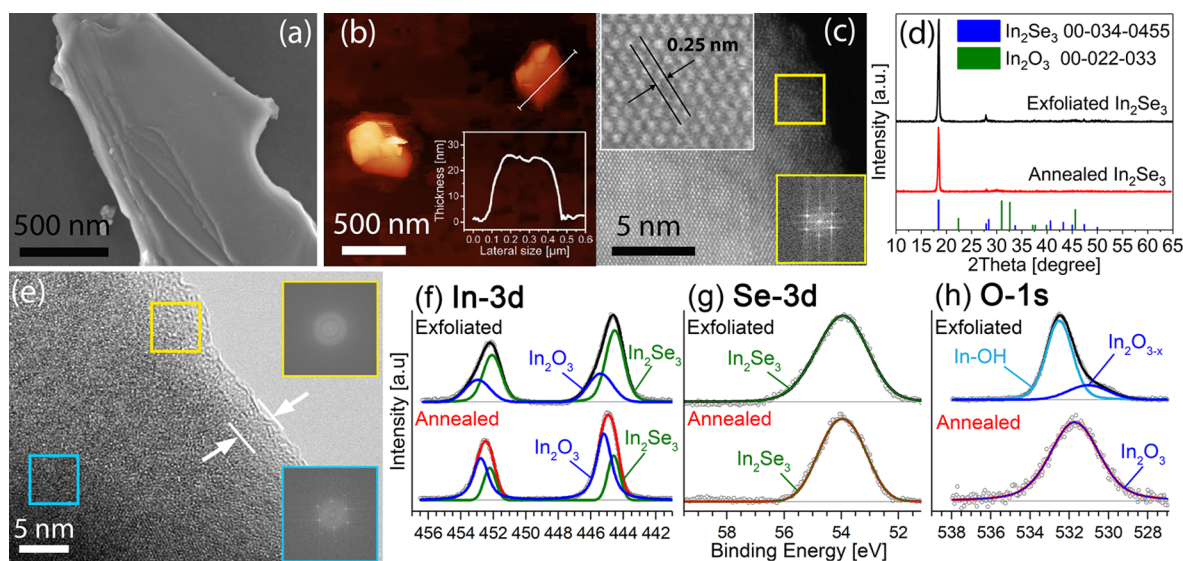
It is reported that, to exfoliate vdW solids, the enthalpy of mixing between the nanosheet's solid surface and the solvent must be minimized, predicting that the surface energy of the nanosheet ( $E_{NS}$ ) be comparable to the total (polar and dispersive) surface energy of the solvent ( $E_S$ ).<sup>24</sup> Considering now the total solvent's surface energy as

$$E_S = (\sigma_S + TS_{surf})$$

where  $\sigma_S$  is the solvent's surface tension and  $S_{surf}$  is the surface entropy ( $\sim 0.1$   $m[N/(m\ K)]$  for liquids); setting  $TS_{surf} \sim 30$   $m[J/m^2]$  for almost all liquids at 20  $^\circ$ C,<sup>25</sup> it turns out that solvent's surface energies  $E_S$  are  $\sim 30$   $m[J/m^2]$  higher than their surface tension  $\sigma_S$ . Consequently, the surface tension  $\sigma_S$  shown in Table 1 must be increased by a factor of 30  $m[J/m^2]$  before matching  $E_S$  with  $E_{NS}$ . On this account, PolarClean, NMP, and IPA, with  $E_S \sim 70$   $m[J/m^2]$ , which is approximately the same order as that of  $WS_2$  nanosheets  $E_{NS} \sim 75$   $m[J/m^2]$ ,<sup>26</sup> possibly explain their superior capability to exfoliate  $WS_2$ , with respect to the others. Conversely, RPDE and urea with associated surface energies of 97 and 104  $m[J/m^2]$  are more effective in exfoliating  $In_2Se_3$  than  $WS_2$ . A tentative explanation, although not entirely satisfactory, could be that  $In_2Se_3$  nanosheets' surface energy  $E_{NS}$  is slightly higher than those reported for TMDs (70–80  $m[J/m^2]$ ), therefore much closer to the surface energy of RPDE and urea solvents than to the others. Data on the surface energy of  $In_2Se_3$  are quite scarce. Proxies to estimate the unknown  $In_2Se_3$  surface energy value could be tentatively those theoretically computed for GaSe and GeSe group III–IV mono-selenides, yielding  $\sim 145$  and  $\sim 220$   $m[J/m^2]$ , respectively.<sup>27</sup> These values are slightly higher than those predicted for 2D TMDs materials, possibly supporting the



**Figure 2.** TG (black line) and DTA (red line) signals of as-exfoliated  $\text{In}_2\text{Se}_3$  flakes performed in air (a) and nitrogen (b) atmospheres in the 25–1050 °C temperature range, with a heating rate of 5 °C/min.



**Figure 3.** (a) SEM image of few flakes of exfoliated  $\text{In}_2\text{Se}_3$ ; (b) AFM picture of exfoliated  $\text{In}_2\text{Se}_3$  flakes with an associated thickness profile; (c) HRTEM image of exfoliated  $\text{In}_2\text{Se}_3$  flakes. Top left inset exhibiting an interplane distance of 0.25 nm corresponding to the (006) interplane distance of  $\text{In}_2\text{Se}_3$ ; (d) grazing-incidence XRD pattern of exfoliated and annealed (at 180 °C for 24 h) samples; (e) HRTEM image of annealed  $\text{In}_2\text{Se}_3$  flakes exhibiting an  $\alpha\text{-In}_2\text{O}_3$  amorphous layer at the edge (white arrows). SAED patterns collected at the edge (yellow) and corresponding to the center of the flake (blue); (f–h) XPS spectra of as-exfoliated (top) and annealed (bottom) flakes: (f) In 3d, (g) Se 3d, and (h) O 1s core-level spectra.

better performances of RPDE and urea solvents to exfoliate  $\text{In}_2\text{Se}_3$ .

### 3.2. Synthesis of the $\alpha\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$ Heterojunction.

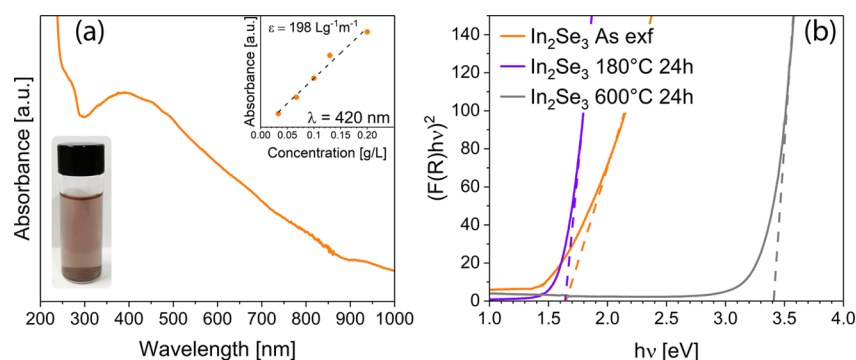
The thermal oxidation properties of  $\text{In}_2\text{Se}_3$  exfoliated flakes, annealed at 5 °C/min in static air and  $\text{N}_2$  from 25 to 1050 °C, have been recorded by TG-DTA simultaneous analysis and shown in Figure 2a,b. No significant weight changes are recorded up to 450 °C temperature in air and in  $\text{N}_2$ , respectively. DTA signals exhibit a modest endothermic peak at 215 °C, ascribed to the transition from  $\alpha\text{-In}_2\text{Se}_3$  to  $\beta\text{-In}_2\text{Se}_3$ .<sup>21,28</sup>

Overall, the TG analysis in air shows a total weight loss of  $-39.5 \pm 0.6\%$ , starting at 450 °C, topping a maximum rate at 575 °C, as indicated by the exothermic DTA peak in Figure 2a. Remarkably, the recorded weight loss of  $-39.5 \pm 0.6\%$  is significantly close to the theoretical weight loss of  $-40.49\%$  corresponding to the formation of  $\text{In}_2\text{O}_3$ . In the same way as to thermal oxidation in air of 2D  $\text{SnSe}_2$  and 2D  $\text{WS}_2$  yielding  $\text{SnO}_2$  and  $\text{WO}_3$ ,<sup>15,29</sup> the formation of  $\text{In}_2\text{O}_3$  proceeds by the complete substitution of the chalcogen atoms (Se) operated by ambient  $\text{O}_2$  according to a thermodynamic spontaneous

process ( $\Delta G < 0$ ). In  $\text{N}_2$  atmosphere, the weight loss of  $-14.9 \pm 0.6\%$  possibly accounts for the sublimation of Se atoms in a similar way to the mechanism discussed for  $\text{SnSe}_2$ .<sup>15</sup> Finally, the sharp DTA endothermic peak at 880 °C is attributed to the melting of  $\text{In}_2\text{Se}_3$ .<sup>30</sup> Exfoliated  $\alpha\text{-In}_2\text{Se}_3$  flakes have been subjected to a controlled oxidation process at 180 °C in air for 24 h at temperatures below the recrystallization temperature of the native oxide (270–300 °C) to promote the formation an amorphous self-assembled  $\alpha\text{-In}_2\text{O}_3$  skin layer over the 2D- $\text{In}_2\text{Se}_3$  flakes.

Figure 3 shows the microstructural features of the flakes before (Figure 3a–c) and after (Figure 3e) annealing, including the corresponding XRD (Figure 3d) and XPS (Figure 3f–h) spectra. Overall, exfoliated  $\text{In}_2\text{Se}_3$  exhibits a stacked-layer morphology (Figure 3a,b) with lateral sizes larger than 500 nm and thicknesses between 25 and 30 nm (see also Figure 1g,h for flakes' size statistics), corresponding to  $\sim 30$  layers (assuming flake's thickness of 6.79 Å and interspace distances of 2.87 Å<sup>21</sup>).

HRTEM analysis of the as-exfoliated  $\text{In}_2\text{Se}_3$  shown in Figure 3c exhibits a highly ordered crystalline texture extending to the



**Figure 4.** Optical characterization of as-exfoliated and annealed  $\text{In}_2\text{Se}_3$ ; (a) UV-vis absorption spectrum of as-exfoliated  $\text{In}_2\text{Se}_3$ . Inset: absorbance coefficient estimated at 420 nm; (b) Tauc plot and BG determination of as-exfoliated and annealed  $\text{In}_2\text{Se}_3$ .

edge of the flake, as confirmed by SAED patterns (yellow square area of the figure). An interplane distance of 0.25 nm (top left inset) corresponding to In atoms tetrahedrally coordinated in a rhombohedral ( $R3m$ ) symmetry of the  $\alpha$ - $\text{In}_2\text{Se}_3$  phase<sup>28</sup> is clearly visible. Such a crystalline structure is confirmed by XRD characterization shown in Figure 3d (black plot), exhibiting, at  $2\theta = 18.4^\circ$ , the characteristic (006) plane orientation of  $\alpha$ - $\text{In}_2\text{Se}_3$  (JCPDS card 00-034-0455).

After annealing at 180 °C for 24 h (Figure 3e), SAED patterns (yellow square) exhibit the formation of an amorphous boundary layer structure approximately 5–7 nm thick, extending from the edge to the center of the flake. The formation of an amorphous  $a$ - $\text{In}_2\text{O}_3$  skin, growing over the underlying 2D  $\text{In}_2\text{Se}_3$  flake, is attested (see the electronic magnification of Figure 3e) by the existence of a disordered structure self-assembled over the whole flake. Remarkably, SAED patterns (blue square) corresponding to the center of the flake confirm the crystalline reflections of the underlying  $\text{In}_2\text{Se}_3$ , attesting the formation of the  $a$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  amorphous/crystalline heterostructure. Additionally, XRD patterns of the annealed flakes (red plot of Figure 3d) highlight a slight reduction of the (006) peak's area of  $\text{In}_2\text{Se}_3$  because of the proceeding of the amorphization process and the absence of new crystalline phases.

Figure 3f–h compares the XPS features of the as-exfoliated and annealed films. The deconvolution of the In 3d core-level spectrum of the exfoliated sample exhibits two contributions having their 5/2 components peaked at the binding energies (BEs) of 444.7 eV (green line) and 445.6 eV (blue line), which can be ascribed to  $\text{In}_2\text{Se}_3$ <sup>31</sup> and  $\text{In}_2\text{O}_3$ <sup>32</sup> phases, respectively, with the former covering 75% of the total spectral area and the latter associated with the remaining 25%. After annealing at 180 °C for 24 h, the opposite takes place. The  $\text{In}_2\text{O}_3$  oxide signal covers 73% of the spectral area, while  $\text{In}_2\text{Se}_3$  has a residual of 27%. Unsurprisingly, considering that XPS information comes from photoelectrons escaping a maximum of up to 10 nm, the occurrence in the In 3d spectra of both signals of  $\text{In}_2\text{O}_3$  and  $\text{In}_2\text{Se}_3$  is further evidence that the  $a$ - $\text{In}_2\text{O}_3$  oxide thickness is limited within 5–7 nm as previously determined by HRTEM characterization (Figure 3e). Regarding the Se 3d spectra, only a single contribution is found in both samples, with its 5/2 component centered at a BE of 54.1 eV corresponding to Se–In bonds in the  $\text{In}_2\text{Se}_3$  structure.<sup>31</sup> Remarkably, the absence of any features at higher binding energies, specifically at  $\sim 58.6$  eV attributed in a previous paper<sup>31</sup> to Se–O bonds, rules out the formation of Se-oxidized phases, indicating that all the Se-bonds belong to  $\text{In}_2\text{Se}_3$ .

Moreover, the formation of an oxide layer over the annealed sample with respect to the exfoliated sample accounts for an attenuation of the Se 3d signal's area of  $\sim 30\%$ . Finally, the O 1s core-level spectra of the as-exfoliated sample (Figure 3h) can be deconvoluted in two different contributions. The lower BE of 531.2 eV, related to defective  $\text{In}_2\text{O}_3$ , namely,  $\text{In}_2\text{O}_{3-x}$  (where  $0 \leq x \leq 1$ ), corresponds to oxygen defects in the matrix of indium oxide,<sup>33</sup> the higher BE, peaking at 532.5 eV, is compatible with hydroxide-related species, as recently described for In-based films.<sup>34</sup> However, after annealing, a clear single contribution centered at 531.9 eV, associated with  $\text{In}_2\text{O}_3$ , excludes the formation of any mixed  $\text{In}_2\text{Se}_{3-3x}\text{O}_{3x}$  ( $0 \leq x \leq 3$ ) oxide phase as previously reported.<sup>31</sup> It may be concluded that both microstructural characterization and chemical composition confirm that annealing at 180 °C for 24 h successfully induces the formation of 5–7 nm thick, amorphous  $a$ - $\text{In}_2\text{O}_3$  skin-layer oxide. The  $a$ - $\text{In}_2\text{O}_3/\text{In}_2\text{Se}_3$  amorphous/crystalline heterostructure passivates the underlying 2D- $\text{In}_2\text{Se}_3$  from further degradation, as will be demonstrated by electrical characterization.

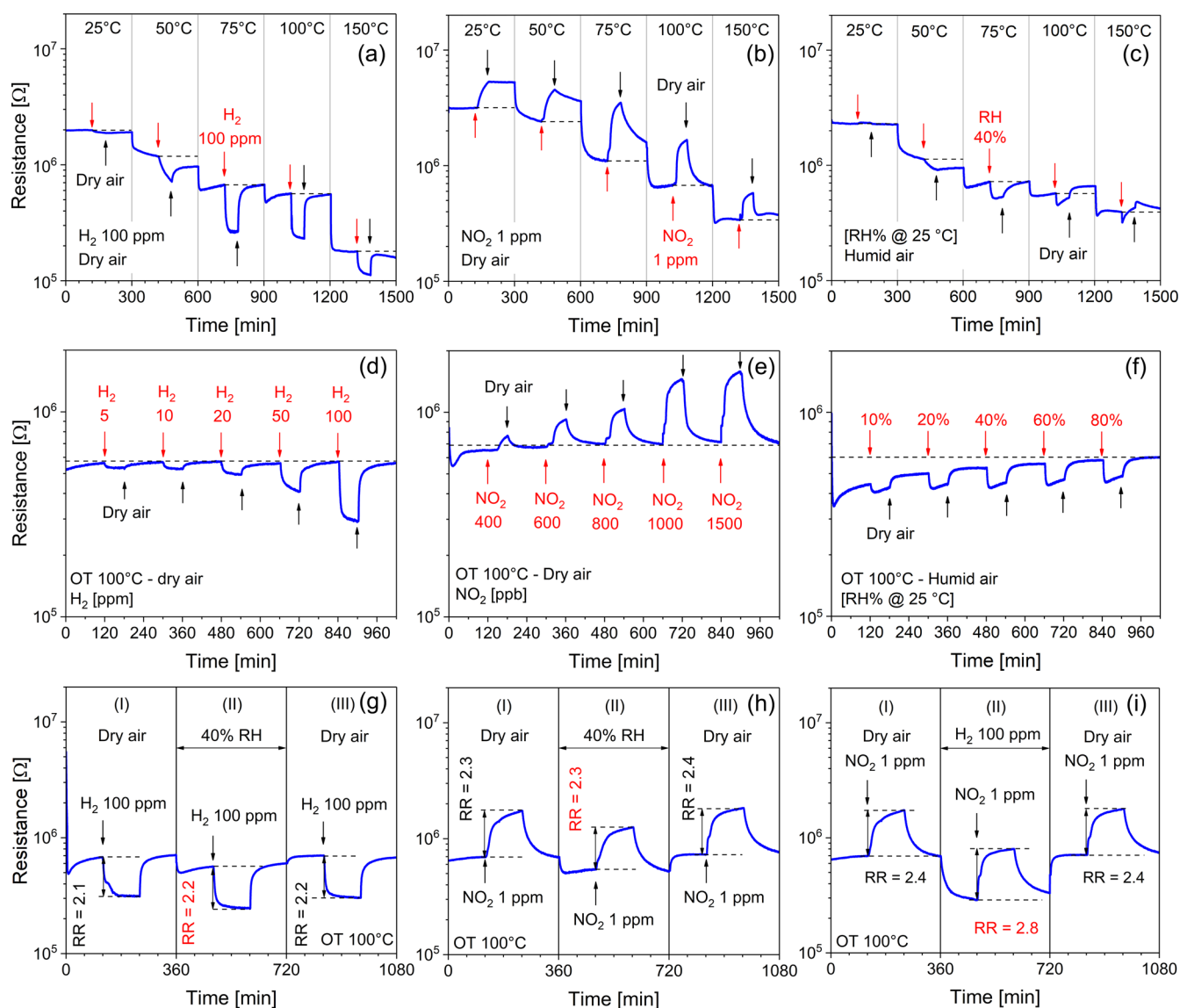
### 3.3. Electrical Properties and Gas Sensing Response.

2D few-layered  $\text{In}_2\text{Se}_3$  has attracted much attention due to its excellent properties and tunable band gaps (BGs) depending on the number of the layers, ranging from  $\sim 1.5$  to 1.7 eV for few layers<sup>35</sup> to 2.27 eV for thin-film photoanodes.<sup>36</sup> Optical BG measurements of as-exfoliated and samples annealed at 180 and 600 °C have been carried out by DRS applying the Kubelka–Munk theory. UV-vis characterization has been performed on ethanol dispersion of as-exfoliated  $\text{In}_2\text{Se}_3$  flakes (0.2 g/L) as shown in Figure 4a.

The as-exfoliated sample broadly absorbs in the whole UV-vis region, with an absorption band edge placed around 850 nm, consistent with  $\text{In}_2\text{Se}_3$  mixed nanoparticles–nanorods.<sup>37</sup> Moreover, both the  $\sim 400$  nm peak and the calculated absorption coefficient [i.e., of 198 L/(g·m)] well agree with those reported for NMP-LPE  $\text{In}_2\text{Se}_3$  nanosheets.<sup>38</sup> Transforming the diffuse reflectance spectra into relative absorption plots using the Kubelka–Munk function,<sup>39</sup> samples' BGs are graphically computed from Figure 4b utilizing Tauc's equation<sup>40</sup>

$$(F(R)hv)^{1/\gamma} = B(h\nu - E_g)$$

where  $h$  is Planck's constant,  $\nu$  is the incident photon energy,  $E_g$  is the optical BG energy, and  $B$  is a constant. Assuming a  $\gamma$  factor of 1/2 corresponding to a direct BG transition for  $\text{In}_2\text{Se}_3$  nanosheets, the extrapolation of Tauc's plots, calculated over five different samples, yields  $1.65 \pm 0.02$ ,  $1.64 \pm 0.03$ , and 3.41



**Figure 5.** Electrical responses in dry air of  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  thin films at different OTs (from 25 to 150 °C) to (a) 100 ppm  $\text{H}_2$ , (b) 1 ppm  $\text{NO}_2$ , and (c) 40% humidity (RH @ 25 °C); (d–f) dynamic electrical responses in dry air at 100 °C OT to (d)  $\text{H}_2$  (5–100 ppm), (e)  $\text{NO}_2$  (400 ppb–1.5 ppm), and (f)  $\text{H}_2\text{O}$  in the range of 10–80% RH (RH at 25 °C). (g) 40% RH humidity cross to 100 ppm  $\text{H}_2$ ; (I) first step in dry air and 100 ppm  $\text{H}_2$ ; (II) second step at 40% RH background and 100 ppm  $\text{H}_2$ ; (III) third step, equivalent to (I), to assess short-term repeatability. (h) 40% RH humidity cross to 1 ppm  $\text{NO}_2$ : (I) first step in dry air and 1 ppm  $\text{NO}_2$ ; (II) second step at 40% RH background and 1 ppm  $\text{NO}_2$ ; (III) third step, equivalent to (I). (i) 100 ppm  $\text{H}_2$  cross to 1 ppm  $\text{NO}_2$ : (I) first step in dry air and 1 ppm  $\text{NO}_2$ ; (II) second step at 100 ppm  $\text{H}_2$  background and 1 ppm  $\text{NO}_2$ ; (III) third step, equivalent to (I).

$\pm 0.03$  eV of the as-exfoliated and 180- and 600 °C-annealed  $\text{In}_2\text{Se}_3$ , respectively. The  $\sim 1.65$  eV BG of the as-exfoliated  $\text{In}_2\text{Se}_3$  well fits with reported values of exfoliated 2D  $\text{In}_2\text{Se}_3$  flakes, with thicknesses ranging from 3 to 28 nm.<sup>35</sup> Most strikingly, the calculated BG of the  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  heterojunction, corresponding to the 180 °C-annealed sample, is  $\sim 1.64$  eV, which is approximately the same of the exfoliated  $\text{In}_2\text{Se}_3$ . Data about BG values of amorphous  $a\text{-In}_2\text{O}_3$  are quite scarce, making the discussion not this straightforward. However, the occurrence of defects in the amorphous structure of  $a\text{-In}_2\text{O}_3$  could entail the formation of intermediate states within the  $a\text{-In}_2\text{Se}_3$  electronic structure, eventually tuning the BG of  $a\text{-In}_2\text{O}_3$  much closer to that of  $\text{In}_2\text{Se}_3$  than crystalline  $\text{In}_2\text{O}_3$ .<sup>31</sup> Another possible explanation could be that the as-grown  $a\text{-In}_2\text{O}_3$  skin with an associated thickness of 5–7 nm is too thin to significantly alter the optical properties of the

sample. Finally, the  $\sim 3.41$  eV BG of the 600 °C-annealed flakes is congruent with the value of  $\sim 3.45$  eV, reported for sintered crystalline  $\text{In}_2\text{O}_3$  thin films prepared by spray pyrolysis.<sup>41</sup>

Electrical tests of the  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  flakes have been carried out in dry/wet air background [0–80% relative humidity (RH) at 25 °C],  $\text{H}_2$  (5–100 ppm), and  $\text{NO}_2$  (0.4–1.5 ppm) at 100 °C OT. Figure 5a–c displays the electrical responses of  $a\text{-In}_2\text{O}_3/\text{In}_2\text{Se}_3$  thin films in dry air to 100 ppm  $\text{H}_2$ , 1 ppm  $\text{NO}_2$ , and 40% RH, respectively, in the OT range of 25–150 °C. The best OT at 100 °C has been selected as a trade-off between the need to regain the BLR (the resistance in air at equilibrium after gas desorption—dotted black lines of Figure 5a–c) and the requisite to yield satisfactory sensor signal variations (i.e.,  $R_a/R_g$  and  $R_g/R_a$ ). Increasing the temperature improves the recovery of the BLR, while reducing the temperature causes

irreversible adsorption of gas molecules, as frequently experienced in MO and 2D layered sensors.<sup>42</sup> Best OTs are found at 75 and 100 °C for H<sub>2</sub> and NO<sub>2</sub>, respectively. The lower hydrogen's OT appears to be related to the smaller electron affinity of H<sub>2</sub> (0.18 eV) with respect to NO<sub>2</sub> (2.3 eV), which facilitates faster desorption. Moreover, sensor's electrical resistances decrease/increase upon exposing it to H<sub>2</sub>/NO<sub>2</sub> gases, confirming an *n*-type response mechanism of the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterojunction. Accordingly, water vapor behaves as a reducing gas, with the best OT at 75 °C (Figure 5c). Overall, the passivation thermal treatment of In<sub>2</sub>Se<sub>3</sub> flakes at 180 °C stabilizes the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterojunction when operating at lower temperatures (75–100 °C), providing, at the same time, an effective solution for fast adsorption/desorption, excellent BLR recovery, and improved sensor's signal reproducibility.

Dynamic gas responses in the dry-air background to H<sub>2</sub>, NO<sub>2</sub>, and different RHs shown in Figure 5d–f exhibit a decrease/increase of the electrical resistance with increasing H<sub>2</sub>/NO<sub>2</sub> concentrations and excellent baseline recovery. In dry air carriers, the experimental limit of detection (LOD) to H<sub>2</sub> is 5 ppm, one of the smallest ever reported for 2D exfoliated TMD and MC interfaces (see Table 2 for comparisons). That to NO<sub>2</sub> is even smaller (400 ppb) considering the stronger attitude of NO<sub>2</sub> to withdraw electrons from the surface. The sensor signal, defined as the ratio  $R_a/R_g$  or  $R_g/R_a$  depending on the reducing/oxidizing nature of the gas, is  $2.1 \pm 0.2$  to 100 ppm H<sub>2</sub> and  $2.3 \pm 0.2$  to 1 ppm NO<sub>2</sub>, providing a  $\pm 0.2$  uncertainty calculated over a set of five identical measure-

ments. These results are comparable to H<sub>2</sub> and NO<sub>2</sub> responses of *a*-SnO<sub>2</sub>/SnSe<sub>2</sub> heterostructures and fully amorphous *a*-SnO<sub>2</sub> 2D interfaces, as previously reported.<sup>14,15</sup>

Cross-sensitivity tests, aimed to assess the variation of the sensor's signal to H<sub>2</sub> or NO<sub>2</sub> target gases when humidity is the interfering one, are shown in Figure 5g,h. According to a standardized procedure, in the first step (I), the sensor is exposed to a dry-air background and the target gases (100 ppm H<sub>2</sub> or 1 ppm NO<sub>2</sub>); in the second step (II), it is exposed to a 40% RH background and target gases; in the third step (III), the sensor is exposed to the same sequence as to (I) to assess the short-term repeatability. According to Figure 5g,h, the BLR always decreases when the sensor is exposed to humidity. Remarkably, sensor's signals (i.e., RRs) to both H<sub>2</sub> and NO<sub>2</sub> gases are almost the same in dry air and 40% RH backgrounds. Specifically, sensor's signal yields  $RR = 2.1 \pm 0.2$  and  $RR = 2.3 \pm 0.2$  to 100 ppm H<sub>2</sub> and 1 ppm NO<sub>2</sub>, respectively, regardless of the humidity content (RH at 40%), apparently suggesting no cross-interference of water vapor to H<sub>2</sub> and NO<sub>2</sub> sensing. In the second run, shown in Figure 5i, the H<sub>2</sub> cross-interference to the NO<sub>2</sub> target gas is displayed. Specifically, the sensor signal to 1 ppm NO<sub>2</sub> significantly increases from  $RR = 2.4 \pm 0.2$  in dry air to  $RR = 2.8 \pm 0.2$  in the 100 ppm H<sub>2</sub> background. Finally, *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> exhibits excellent short-term repeatability properties, as attested by comparing the electrical responses of panel (I) and panel (III) for all the investigated sequences. The selectivity response, obtained by comparing the sensor's RRs to 100 ppm H<sub>2</sub> and 1 ppm NO<sub>2</sub> in dry air, with respect to different gases and organic vapors, is also shown in Figure S6. The humidity cross-interference to 100 ppm H<sub>2</sub> and 1 ppm NO<sub>2</sub> target gases in the whole 0–80% RH range is displayed in Figure 6a,b. Remarkably, sensor's signal response to 100 ppm H<sub>2</sub> is the same in the humidity range of 0–60% while that to NO<sub>2</sub> remains constant in a narrower range (0–40%). These results confirm, as previously found by cross-test investigations (Figure 5g,h), that no cross-sensitivity effects of water vapor are detectable in the 0–40% range for both gases.

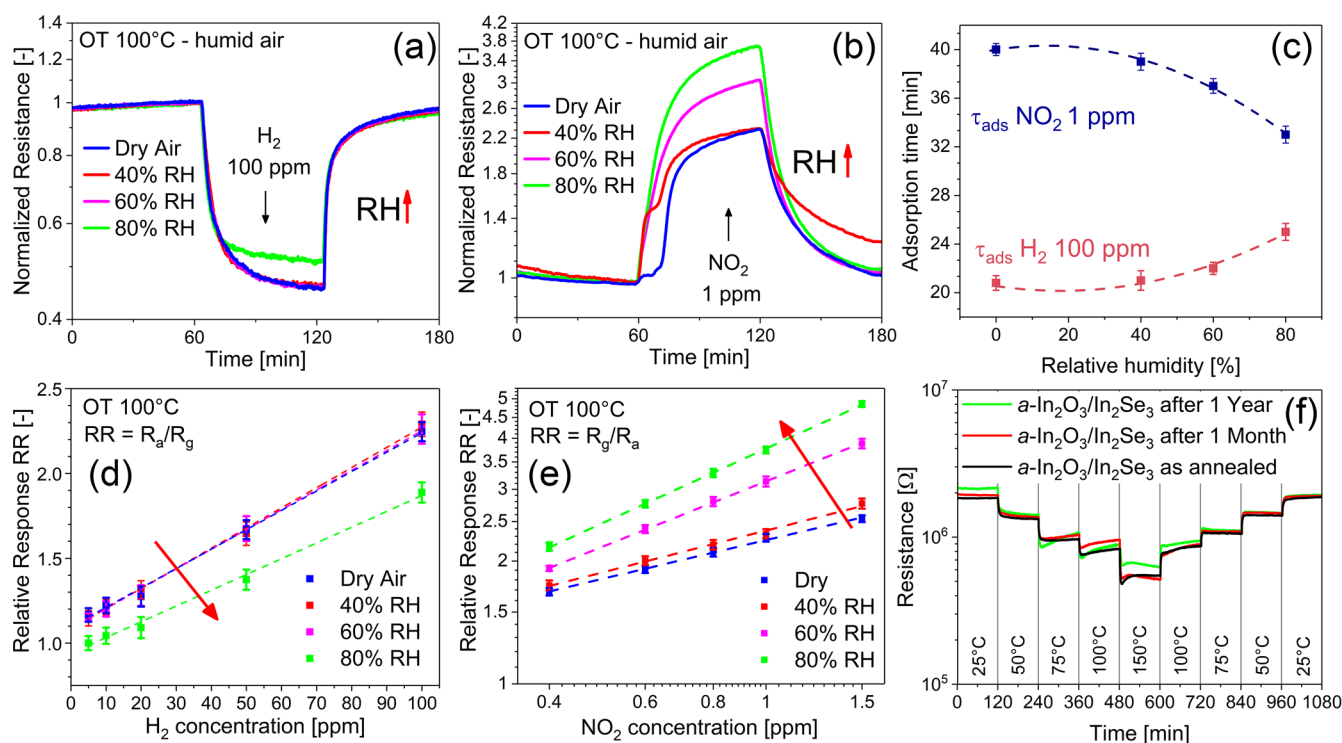
Notably, with increasing RH (as to the direction of the red arrow of Figure 6a,b), water vapor inhibits/enhances the H<sub>2</sub>/NO<sub>2</sub> sensor's signal, displaying an antisnergistic/synergistic effect upon sensor's signal amplitude. Specifically, signal responses to 100 ppm H<sub>2</sub> decrease from  $2.1 \pm 0.2$  to  $1.7 \pm 0.2$  from 0 to 80% RH. Conversely, that to 1 ppm NO<sub>2</sub> increases from  $2.4 \pm 0.2$  to  $3.8 \pm 0.2$ . The log–log calibration plots of the sensor's signal (i.e.,  $R_a/R_g$  for H<sub>2</sub> and  $R_g/R_a$  for NO<sub>2</sub>) in the whole gas concentration range and different RH values are shown in Figure 6d,e. At RH % higher than 60 and 40%, the calibrating H<sub>2</sub> and NO<sub>2</sub> lines shift downward/upward (as to the direction of the red arrows of Figure 6d,e), respectively. Congruent with this tendency, water vapor has a remarkable effect to change the LOD. Specifically, the decrease/increase of the H<sub>2</sub>/NO<sub>2</sub> slopes of the calibrating lines at 80% RH (with respect to that at 0% RH) downshifts/upshifts the theoretical LOD (obtained by graphically extrapolating the calibrating lines to the intercept drawn at  $RR = 1$ ). Most strikingly, a high level of humidity has a strong influence on the adsorption time. Increasing the RH increases/decreases the adsorption times to H<sub>2</sub> and NO<sub>2</sub> as shown in Figure 6c.

Combining these observations, it may be concluded that depending on the reducing/oxidizing nature of the interfering gas, provided an *n*-type response of the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub>

**Table 2. Comparison of the H<sub>2</sub> Gas Sensing Performances of Different Sensor's Interfaces Obtained by Normalizing Literature Data to 10 ppm H<sub>2</sub> in Dry Air**

Sensing material and structure	LOD [ppm] <sup>a</sup>	RRs 10 ppm H <sub>2</sub> <sup>b</sup>	OT [°C]	refs
MOS–Noble Metal/MOS–MOS Heterostructures				
In <sub>2</sub> O <sub>3</sub> mesoporous	0.1	1.8	260	43
Au–In <sub>2</sub> O <sub>3</sub> core–shell	2	5.0	300	44
Au–SnO <sub>2</sub> nanoparticles	5	6.2	250	45
Pd–SnO <sub>2</sub> nanofibers	0.25	12.6	160	46
CeO <sub>2</sub> –In <sub>2</sub> O <sub>3</sub> hollow sphere	1	12.2	160	47
NiO–ZnO nanofibers	5	1.5	200	48
2D-Graphene (G-GO-rGO)–NP-Decorated/G-rGO				
G-decorated Pd NCL	0.1	1.06	25	49
rGO-decorated Pt NPs	1	1.04	25	50
G-decorated SnO <sub>2</sub> NPs	1	2.00	50	51
2D TMDs–MCs.–NP-Decorated/TMDs–MCs				
Pt NPs–MoS <sub>2</sub>	10	1.20	150	52
2D–MoS <sub>2</sub>	10	1.10	25	53
Pd–SnO <sub>2</sub> /MoS <sub>2</sub>	30	1.02 (@100 ppm)	25	54
Pd NPs–MoS <sub>2</sub>	50	1.18 (@100 ppm)	25	55
2D TMDs–MCs–Heterostructures				
<i>a</i> -WO <sub>3</sub> /WS <sub>2</sub> amorphous/cryst	5	2.80	150	29
<i>a</i> -SnO <sub>2</sub> /SnSe <sub>2</sub> amorphous/cryst	5	1.22	100	14
<i>a</i> -In <sub>2</sub> O <sub>3</sub> /In <sub>2</sub> Se <sub>3</sub> amorphous/cryst	5	1.20	100	this work

<sup>a</sup>Experimental LOD. <sup>b</sup>RR =  $R_a/R_g$  for *n*-type and  $R_g/R_a$  for *p*-type sensors.



**Figure 6.** (a,b) Sensor's responses to increasing RH % content (as to the direction of the red arrow) to 100 ppm H<sub>2</sub> and 1 ppm NO<sub>2</sub>; (c) adsorption times to 100 and 1 ppm NO<sub>2</sub> as a function of the RH % content; (d,e) log/log calibration plots at 100 °C OT and increasing RH % contents (as to the red arrow) from 0 to 80% (RH at 25 °C) to H<sub>2</sub> and NO<sub>2</sub>; (f) BLR variations in dry air at different OTs by increasing/decreasing the OT in the 25–150–25 °C range. Black line: as-annealed *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub>; red: measured after 1 month; green: after 1 year of continuous delivery at 100 °C OT. Associated standard deviations  $\sigma \pm 0.2$ , calculated over a set of five identical measurements.

heterojunction, the increase of the RH % content displays an antisnergistic/synergistic effect upon (i) H<sub>2</sub>/NO<sub>2</sub> sensor signal variation and (ii) H<sub>2</sub>/NO<sub>2</sub> adsorption times. Overall, the single effect of adsorbing water is to decrease the absolute value of the BLR of the material for all the investigated combinations. Finally, the long-term stability properties of the BLR and sensor's signal (RRs) have been recorded by comparing the electrical signals of the as-annealed *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> with that of the same sample after 1 month and 1 year of continuous delivery at a 100 °C OT. Notably, excellent reproducibility of the BLR at different OTs between 25 and 150–25 °C is displayed in the dry-air background in Figure 6f. Moreover, the excellent long-term stability for 1 year of both BLR and sensor signals to H<sub>2</sub> and NO<sub>2</sub>, at 75 and 100 °C OTs as shown in Figure S7, presents indirect proof that the oxide skin layer effectively passivates the underlying 2D-In<sub>2</sub>Se<sub>3</sub> from spontaneous degradation.

Normalizing from literature, we finally compared the experimental LOD, sensor's signal responses (RRs at 10 ppm H<sub>2</sub>), and OTs of best-performing materials and structures in Table 2. MO undoubtedly perform better but at the expense of higher OTs. Surprisingly, a very limited number of 2D layered interfaces comprising 2D carbon materials and TMDs/MCs have been found to respond to few ppm H<sub>2</sub> with low LODs. Amorphous/crystalline *a*-MO<sub>x</sub>/TMD heterostructures operating at mild OTs (100 °C) show better features as compared to simple 2D TMD/MC interfaces, possibly on account of the role played by the oxide amorphous layer to enhance gas response.

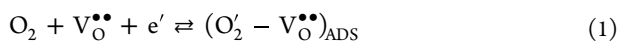
**3.4. Discussion.** We demonstrated (Figure 1f) that both exfoliation yields and ARs (i.e., the ratio of lateral dimension to

the thickness) (Figure 1g,h) depend on the physio-chemical features of the solvent and the sonication/centrifugation process conditions. However, to date, there are no literature investigations demonstrating possible correlations between the AR and the gas sensing response. Conversely, as recently reported, flakes' edges with high concentrations of dangling bonds with respect to terraces, represent the most reacting parts of the nanosheets.<sup>56</sup> These findings address the usefulness of maximizing the length of the exposed edges, with respect to the surface, which enhances the chemical activity of the gas molecules. The passivation strategy to preserve TMDs/MCs from environmental degradation, accounts for controlled oxidation in air of exfoliated In<sub>2</sub>Se<sub>3</sub>, yielding the formation of an amorphous skin, which guarantees excellent baseline and sensor's signal reproducibility after 1 year of delivery as shown in Figure 6f. Overall, the oxidation mechanism induces an isomorphic conversion of 2D In<sub>2</sub>Se<sub>3</sub> into 2D *a*-In<sub>2</sub>O<sub>3</sub>, thus preserving in the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> annealed flakes all the surface-to-volume advantages of the precursor 2D In<sub>2</sub>Se<sub>3</sub> layered structure.

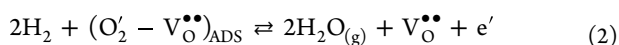
The oxidation process of TMDs (MX<sub>2</sub>: M = Mo, W, X = S, Se) and MCs (MX: M = Ga, In, Ge, Sn; X = S, Se) in air consists of: (i) the formation of edge and/or basal plane vacancy (i.e., either metal or chalcogen), (ii) the diffusion of molecular oxygen (O<sub>2</sub>) inside the material, (iii) the dissociative chemisorption of O<sub>2</sub> as atomic oxygen (O) over the vacancy, and (iv) the annihilation of the vacancy by the formation of metal–oxygen (M–O) bonds.<sup>57</sup> Specifically, it has been theoretically found that the formation of a single selenium vacancy is energetically favorable in the InSe monolayer (~1.9 eV) and that the O<sub>2</sub> dissociation over the Se vacancy is highly

exothermic (4.57 eV),<sup>58</sup> leading to the formation of a stable In–O bond. Besides energetics, from a kinetic point of view, it can be possibly assumed that the oxidation process is governed by the diffusion of molecular oxygen (O<sub>2</sub>) from the surface of the 2D material to the selenium vacancy sites, following a shrinking-core reaction mechanism,<sup>59</sup> as detailed in Section S4. According to this model, the as-formed *a*-In<sub>2</sub>O<sub>3</sub> amorphous skin builds up a diffusion boundary layer which controls the O<sub>2</sub> diffusion toward the unreacted particle core of In<sub>2</sub>Se<sub>3</sub>. It turns out that the oxidation rate of rectangular stacked nanosheets of In<sub>2</sub>Se<sub>3</sub> monolayers depends on the square root of the annealing time and exponentially on the annealing temperature, highlighting the higher significance of annealing temperature with respect to the duration time. Remarkably, the exponential temperature correlation would imply that the oxidation rate of In<sub>2</sub>Se<sub>3</sub> is negligible at room temperature (i.e., referred to as spontaneous oxidation), which is apparently in contrast to the 25% of *a*-In<sub>2</sub>O<sub>3</sub> phase found by XPS analysis, in the as-exfoliated In<sub>2</sub>Se<sub>3</sub> (Figure 3f). Notably, XPS measurements, shown in Figure S4, show that pristine commercial crystals are already oxidized, probably depending on the synthesis route, with approximately the same amount of *a*-In<sub>2</sub>O<sub>3</sub> phase content as the as-exfoliated sample. As a consequence, spontaneous oxidation at room temperature under dry/wet environmental laboratory conditions or due to the exfoliation process can be excluded. However, the occurrence of a preoxidized state (i.e., 25% *a*-In<sub>2</sub>O<sub>3</sub>) in the as-exfoliated samples does not interfere with the thermal oxidation synthesis process at a higher temperature. Moreover, congruent to the kinetic model, by changing the annealing temperature, it is possible to modulate and control the thickness of the skin oxide, although focused experiments are needed to validate the model. In conclusion, annealing at 180 °C for 24 h enhances the formation of an *a*-In<sub>2</sub>O<sub>3</sub> skin layer, which kinetically inhibits oxygen diffusion at lower temperatures, therefore effectively passivating the In<sub>2</sub>Se<sub>3</sub>/In<sub>2</sub>O<sub>3</sub> heterostructure from further oxidation. The stabilization of the sensor's BLR for a period of 1 year, at a 100 °C OT (Figure S7), is indirect proof of the effectiveness of the passivation procedure.

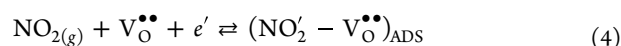
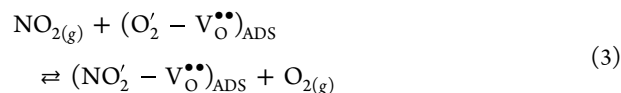
The gas sensing response of the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterostructure behaves like an *n*-type crystalline In<sub>2</sub>O<sub>3</sub> chemoresistive H<sub>2</sub>/NO<sub>2</sub> sensor.<sup>60</sup> The higher sensor signal amplitude and better detection limits of the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterostructure with respect to 2D-TMDs/MCs shown in Table 2 account for the presence of Se and O<sub>2</sub> vacancies within the amorphous structure of *a*-In<sub>2</sub>O<sub>3</sub><sup>31</sup> which largely improve gas adsorption. According to the literature, following the Kröger–Vink notation, sensor signal variation is driven by redox surface reactions between H<sub>2</sub>, NO<sub>2</sub>, and H<sub>2</sub>O molecules with positively ionized oxygen vacancies (V<sub>O</sub><sup>••</sup>), ionosorbed oxygen (O<sub>2</sub><sup>'</sup>), hydroxyls (OH<sup>'</sup>), rooted (OH)<sub>O</sub><sup>•</sup>, and lattice In<sub>in</sub><sup>x</sup> and Se<sub>se</sub><sup>x</sup>. All these species compete at the same time with the interfering gases. Specifically, at low OT (*T* < 130 °C) and dry air, oxygen adsorbs over (V<sub>O</sub><sup>••</sup>), according to reaction (1), increasing the sensor's resistance



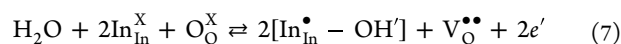
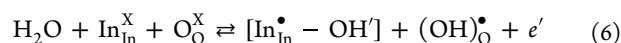
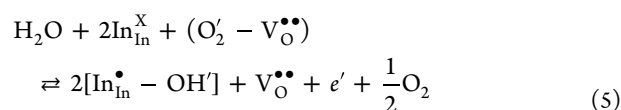
Following a combustion mechanism, hydrogen reacts with ionosorbed oxygen, releasing a vacancy and an electron, decreasing the sensor's resistance as per reaction (2).



NO<sub>2</sub> gas, conversely, can either react with (O<sub>2</sub><sup>'</sup> – V<sub>O</sub><sup>••</sup>) according to reaction (3) (in this case, NO<sub>2</sub>, being more electronegative simply displaces adsorbed oxygen, broadening the depletion layer) or directly react with free oxygen vacancies according to reaction (4), forming nitrite species and consuming electrons. In both cases, NO<sub>2</sub> gas increases the sensor's resistance.

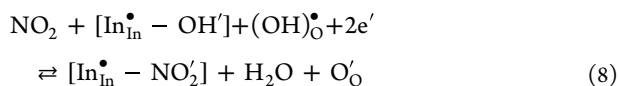


Water vapor, on the other hand, reacts with sensor's surface mostly in a dissociative chemisorption mode,<sup>15</sup> releasing electrons in the conduction band of the material, according to reactions (5–7)



In the first mechanism (5) proposed by Heinrich and Cox,<sup>61</sup> water vapor displaces adsorbed oxygen, forming hydroxyls (OH<sup>'</sup>) and releasing a vacancy (V<sub>O</sub><sup>••</sup>) and e'. This mechanism seems to be supported by a strong reduction of the O<sub>2</sub>/OH<sup>'</sup> radicals' ratio in humid conditions, as recorded by EPR measurements for crystalline SnO<sub>2</sub>.<sup>62</sup> In the second mechanism, proposed by Heiland and Kohl<sup>63</sup> (6,7), either rooted (OH)<sub>O</sub><sup>•</sup> or two terminal (OH<sup>'</sup>) groups and an oxygen vacancy are formed. All these mechanisms, producing electrons, are congruent with the decrease of the BLR in 40% RH as recorded in Figure 5g,h. Remarkably, as shown in Figure 6a,d, increasing the humidity content decreases the H<sub>2</sub> sensor's signal amplitude, following the same response as to CO in SnO<sub>2</sub> and In<sub>2</sub>O<sub>3</sub> crystalline MOs.<sup>64–66</sup> Combining these observations, the decrease of the H<sub>2</sub> sensor signal with increasing humidity, accounts first for the synergistic action of reactions (2, 5–7) which release electrons/vacancies (leading to a resistance decrease), followed by the recombination of air oxygen with the as-formed electrons/vacancies, as for reaction (1) (leading to a resistance increase). An equilibrium is established, where a resistance decrease is partially mitigated by a resistance increase. Overall, humidity negatively affects sensor's signal amplitude to H<sub>2</sub>, according to an antisnergistic effect. Conversely, increasing humidity increases the NO<sub>2</sub> sensor signal, as shown in Figure 6b,e. Specifically, in dry air, NO<sub>2</sub> depletes electrons from the conduction band, increasing the resistance according to reaction (4); meanwhile, humid air releases vacancies and electrons, as to reactions (5–7). Considering now that NO<sub>2</sub> gas has a stronger oxidizing activity than oxygen (its electron affinity *E*<sub>ea</sub>(NO<sub>2</sub>) = 2.27 eV is higher than *E*<sub>ea</sub>(O<sub>2</sub>) = 0.44 eV), in NO<sub>2</sub> atmosphere, reaction (4) will prevail with respect to the reoxidation of the surface operated by air-oxygen (1). Combining these observations, the resistance increase due to the NO<sub>2</sub> reaction (4) in dry air (resistance increases) is boosted, in humid air, by further NO<sub>2</sub> reaction with electrons/vacancies released by water adsorption (5–7) (resistance increase). Overall humidity

positively affects sensor's signal increase to NO<sub>2</sub> according to a synergistic effect. Interestingly, operando DRIFTS investigations over crystalline In<sub>2</sub>O<sub>3</sub><sup>67</sup> at 130 °C pointed out that hydroxyl groups formed by water vapor adsorption (5) can be easily displaced by NO<sub>2</sub> adsorption, according to a "wet" adsorption mechanism, as to reaction (8). This mechanism, as previously reported,<sup>61</sup> clearly accounts for the NO<sub>2</sub> gas sensor signal increase with increasing water vapor content, at low temperatures ( $T < 130$  °C)



The increase/decrease of the H<sub>2</sub>/NO<sub>2</sub> adsorption times with increasing humidity content, as shown in Figure 6c, is finally explained considering that H<sub>2</sub> gas reacts with only (O'<sub>2</sub> - V<sup>••</sup><sub>O</sub>) (2), whereas NO<sub>2</sub> reacts with (O'<sub>2</sub> - V<sup>••</sup><sub>O</sub>), V<sup>••</sup><sub>O</sub> and OH' (3, 4, 8). In a H<sub>2</sub> atmosphere, increasing the humidity content decreases the concentration of (O'<sub>2</sub> - V<sup>••</sup><sub>O</sub>) due to the competitive action of H<sub>2</sub> and H<sub>2</sub>O over the same adsorption site [reactions (2) and (5)]. In this case, H<sub>2</sub> adsorption time increases with increasing humidity content. Conversely, the increase of the concentration of V<sup>••</sup><sub>O</sub> and OH' adsorption sites with increasing humidity content [reactions (5,6,7)] increases the number of available adsorption sites for NO<sub>2</sub>, therefore, decreasing adsorption time. It should be noted that, according to a previously reported mechanism,<sup>68</sup> all the above-mentioned reactions tune the extension of the depletion layers (known as the Debye length) within the *a*-In<sub>2</sub>O<sub>3</sub> oxide skin. With an associated *a*-In<sub>2</sub>O<sub>3</sub> skin thickness of 5–7 nm (see Figure 3e), considering a Debye length between 5 and 50 nm, it is reasonable that the *a*-In<sub>2</sub>O<sub>3</sub> oxide skin is fully or partially depleted of electrons when in contact with oxidizing gas [NO<sub>2</sub>—reactions (3, 4, 8)] or reducing [H<sub>2</sub>, H<sub>2</sub>O—reactions (2, 5–7)] gases, respectively. Assuming a parallel contribution scheme of the *a*-In<sub>2</sub>O<sub>3</sub> skin and In<sub>2</sub>Se<sub>3</sub> bulk to the conduction model of the whole flake,<sup>68</sup> taking into consideration that atomically thin layers of In<sub>2</sub>O<sub>3</sub> exhibit 2 orders-of-magnitude enhancement of conductivity and mobility with respect to few-layer In<sub>2</sub>Se<sub>3</sub>,<sup>69</sup> we may conclude that the outer *a*-In<sub>2</sub>O<sub>3</sub> shell governs the conduction mechanism of the whole flake. This is achieved, depending on the nature of the adsorbing gas, by the modulation of the Debye length within the *a*-In<sub>2</sub>O<sub>3</sub> skin. Lastly, flake-to-flake contacts between loosely stacked flakes built up, as previously depicted,<sup>70</sup> Schottky barriers between the flakes, modulating the conductivity similarly to the conduction mechanism of unsintered MO nanoparticles in traditional chemoresistive sensors.

#### 4. CONCLUSIONS

We have reported a simple thermal treatment process in air, carried out at temperatures below the recrystallization temperature of the native oxide that induces the formation of ~5–7 nm thin, amorphous oxide layers over liquid-phase exfoliated (~25–30 nm thick) 2D-In<sub>2</sub>Se<sub>3</sub> flakes, leading to the formation of an amorphous/crystalline *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterostructure. We found that the self-assembled growth of the native oxide leads to an isomorphic conversion of 2D-In<sub>2</sub>Se<sub>3</sub> into *a*-In<sub>2</sub>O<sub>3</sub>, implicating that the layered shape of the precursor 2D-In<sub>2</sub>Se<sub>3</sub> exfoliated is fully maintained. We also found that the *a*-In<sub>2</sub>O<sub>3</sub>/In<sub>2</sub>Se<sub>3</sub> heterostructure responds to H<sub>2</sub> and NO<sub>2</sub> gases with associated detection limits as low as 5 ppm

and RR ( $R_a/R_g$ ) = 1.2 to 10 ppm H<sub>2</sub>, the highest with respect to only exfoliated 2D TMD/MC sensors. Most strikingly, the *a*-In<sub>2</sub>O<sub>3</sub> skin oxide effectively passivates the underlying 2D-In<sub>2</sub>Se<sub>3</sub> from further ambient oxidation, as attested by the excellent long-term reproducibility of the BLR and sensor's signal to H<sub>2</sub> and NO<sub>2</sub>, over 1 year delivery at 100 °C OT. We finally studied a possible gas sensing mechanism which accounts for a synergistic/antisynergistic effect played by humid water over the H<sub>2</sub>/NO<sub>2</sub> sensor signal amplitude and adsorption times. In conclusion, this methodology, which can be successfully extended to a large variety of TMDs and MCs, represents a possible strategy to improve the long-term stability of the electrical properties of TMDs/MCs devices for gas sensing applications.

#### ■ ASSOCIATED CONTENT

##### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsnm.3c00380>.

Experimental section, Rhodiasolv RPDE's FTIR characterization, microstructural characterization, shrinking-core model of In<sub>2</sub>Se<sub>3</sub> oxidation, and gas sensing responses (PDF)

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##### Author Contributions

Author Credit statement: V.P. performed all material synthesis, conducted the research investigation process, evaluated the measurements, and drafted the manuscript. J.D.S. performed material characterization and conducted a formal analysis of the collected data by the application of dedicated software. V.R. performed material synthesis and electrical characterization addressing a model of the oxidation mechanism of In<sub>2</sub>Se<sub>3</sub>. L.L. carried out the research and investigation process by performing XPS experiments. C.C. addressed the goals of the research, supervised and coordinated experiments, and finalized the manuscript.

## Notes

The authors declare no competing financial interest.

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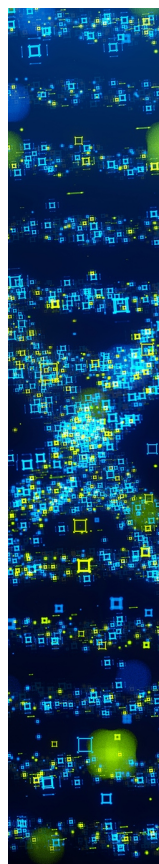
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